



# Temporal Waveguiding in a Dispersive Time-Varying Medium

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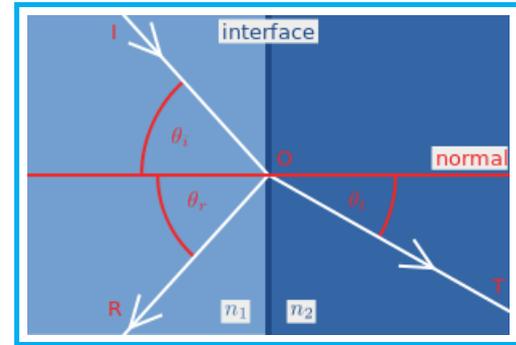
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# Introduction

- Wave propagation in a time-varying medium has a long history.  
F. R. Morgenthaler, IRE Trans. Microwave Theory Tech. **6**, 167 (1958);  
C. Caloz and Z. Deck-Léger, Parts I & II, IEEE Trans. Anten. Propag. **68**, 1569-1598 (2020).
- Simple case: A temporal boundary with different  $\epsilon$  on its two side.  
T. Ruiz et al., IEEE Trans. Anten. Propag. **26**, 358 (1978);  
Y. Xiao, D. N. Maywar, and G. P. Agrawal, Opt. Lett. **39**, 574 (2014).
- A moving temporal boundary in a dispersive medium is of practical interest.  
Plansinis, Donaldson, Agrawal, PRL **115**, 183901 (2015);  
Plansinis, Donaldson, Agrawal, IEEE JQE **52**, 6100708 (2016).
- Analog of total internal reflection can be used for temporal waveguiding  
Plansinis, Donaldson, Agrawal, JOSA B **33**, 1112 (2016).

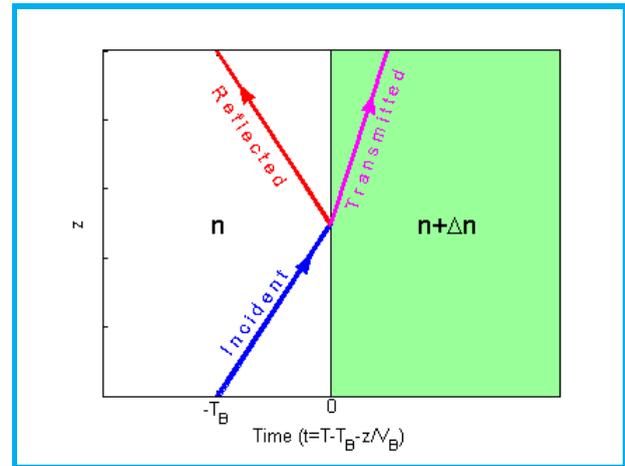
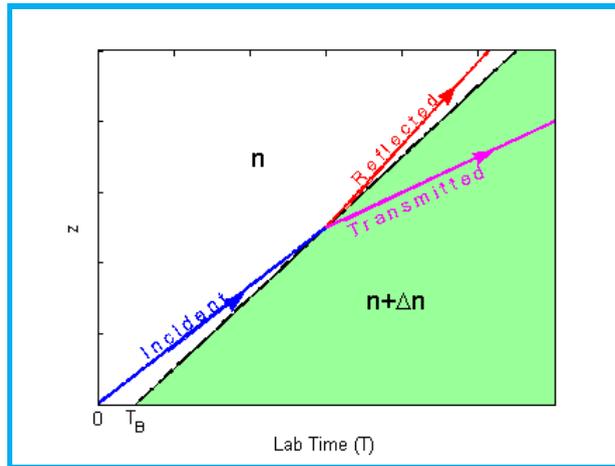
# Temporal Reflection and Refraction

- Reflection and refraction of optical beams at a spatial boundary are well-known phenomena.
- What is the temporal analog of these two optical phenomena?



- What happens when an optical pulse arrives at a temporal boundary across which the refractive index changes from  $n$  to  $n + \Delta n$ ?
- At a spatial boundary, frequency is preserved but momentum can change.
- At a temporal boundary, momentum is preserved but frequency can change.
- It is difficult to change  $n$  at a specific time across the entire medium. A moving temporal boundary is easier to realize.

# Reflection at a Moving Boundary



- Both the pulse and temporal boundary travel forward at different speeds.
- It is convenient to work in a moving frame at which the boundary is stationary ( $t = T - z/V_B$ ).
- Momentum (or the wave vector) is conserved in the moving frame.

# Pulse Propagation in a Dispersive Medium

- Consider an optical pulse propagating inside a medium with the dispersion relation  $\beta(\omega)$ .
- Temporal discontinuity at  $t = T_B$  can be incorporated through a step function. Its use provides the dispersion relation:

$$\beta(\omega) = \beta_0 + \Delta\beta_1(\omega - \omega_0) + \frac{\beta_2}{2}(\omega - \omega_0)^2 + \beta_B H(t - T_B).$$

- $\Delta\beta_1 = \beta_1 - 1/V_B$  is pulse's speed relative to the temporal boundary located at  $t = T_B$ ;  $H(t - T_B)$  is the Heaviside step function.
- $\beta_B = k_0\Delta n$  when refractive index changes by  $\Delta n$  for  $t > T_B$  ( $k_0 = 2\pi/\lambda$ ).
- The dispersion relation can be used to investigate changes in pulse's shape and spectrum occurring when the pulse arrives at the boundary.

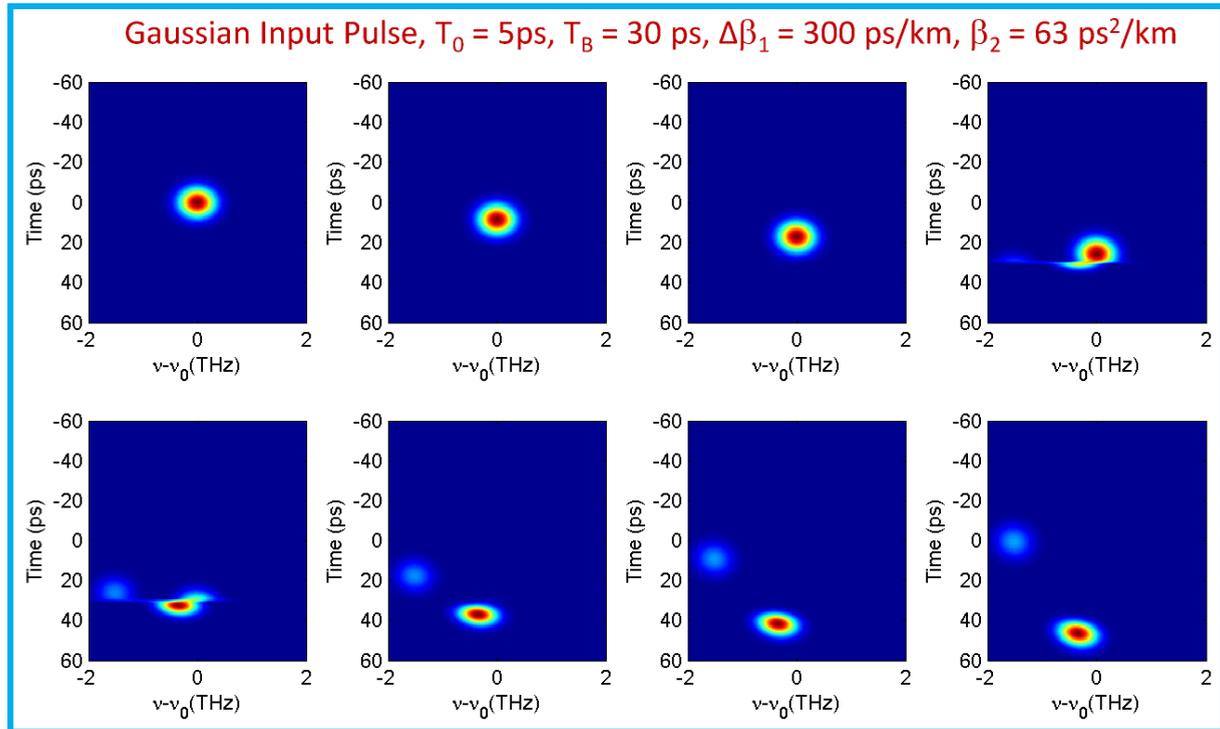
# Equation for the Pulse Envelope

- Slowly varying envelope of the pulse satisfies

$$\frac{\partial A}{\partial z} + \Delta\beta_1 \frac{\partial A}{\partial t} + \frac{i\beta_2}{2} \frac{\partial^2 A}{\partial t^2} = i\beta_B H(t - T_B)A.$$

- This linear equation can be solved with the split-step Fourier method to study what happens at a moving temporal boundary.
- Numerical results show that the pulse splits into two parts with the spectra shifted in frequency by different amounts.
- Two parts propagate at different speeds because of chromatic dispersion.
- Reflected part moves slower and never crosses the temporal boundary.
- Fraction of energy in the reflected pulse depends on the index change  $\Delta n$ . Total reflection is possible for large  $\Delta n$ .

# Simulations for Gaussian pulses

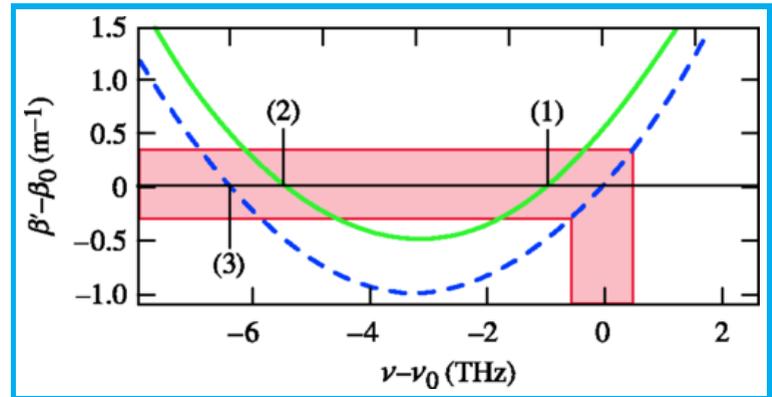


Spectrograms at different distances showing pulse's splitting at the TB.

# Momentum Conservation

Conservation of momentum explains all results:

- Blue curve for  $t < T_B$
- Green curve for  $t > T_B$
- Red region: pulse spectrum



- Reflection corresponds to solution (3) on the blue curve.
- Refraction corresponds to solution (1) on the green curve.
- Solution (2) is not physical because its slope is opposite to that of (1).
- Both reflection and refraction manifest as red-shifted pulses; blue shifts occur if  $\beta_2$  or  $\Delta n$  is negative.

# Spectral Shifts at the Temporal Boundary

- Momentum conservation provides analytical expressions for frequency shifts:

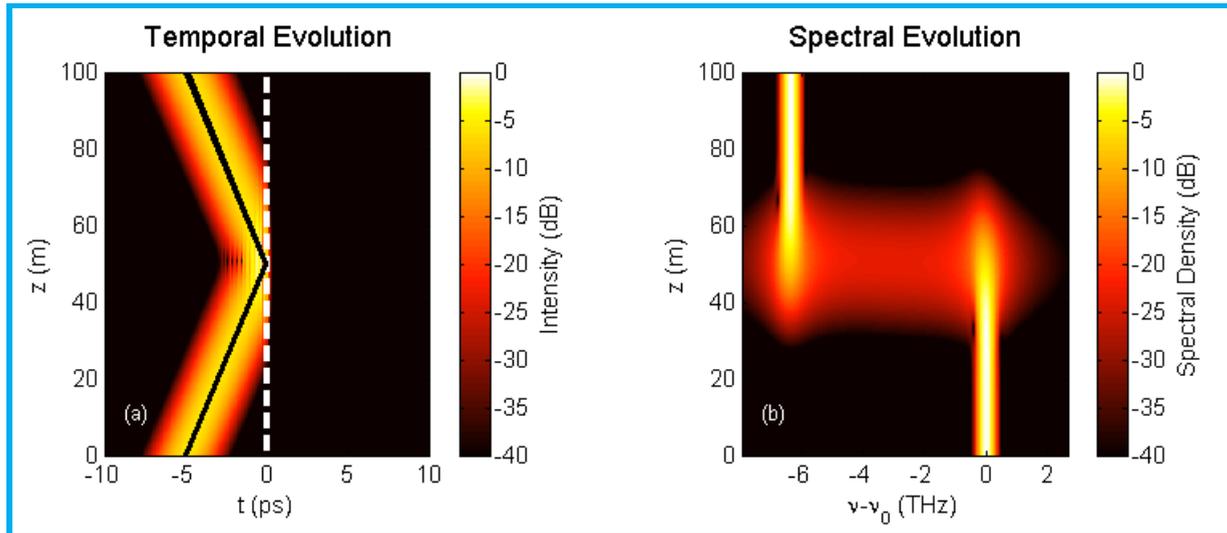
$$\omega_r = \omega_0 - 2(\Delta\beta_1/\beta_2),$$

$$\omega_t = \omega_0 + \frac{\Delta\beta_1}{\beta_2} \left( \sqrt{1 - \frac{2\beta_B\beta_2}{(\Delta\beta_1)^2}} - 1 \right).$$

## Total Internal Reflection

- For large  $\beta_B$ ,  $\omega_t$  becomes a complex quantity. This situation corresponds to temporal analog of total internal reflection (TIR).
- The TIR condition is  $(2\beta_B\beta_2) > (\Delta\beta_1)^2/$ . Temporal TIR occurs only if  $\beta_B$  and  $\beta_2$  have the same signs.
- In the case of  $\beta_2 > 0$ ,  $n$  increases after the temporal boundary. The opposite scenario occurs for  $\beta_2 < 0$ .

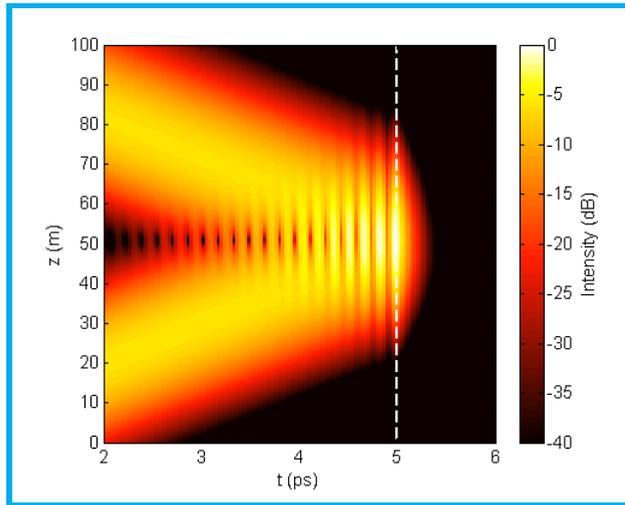
# TIR of Gaussian Pulses



Plansinis, Donaldson, Agrawal, PRL **115**, 183901 (2015).

- Index change was large enough ( $\Delta n \sim 10^{-6}$ ) to satisfy the TIR condition.
- Entire pulse energy gets reflected at the temporal boundary.
- Spectrum shifts by  $> 6$  THz after TIR of the pulse at the boundary.

# TIR of Gaussian Pulses (cont.)

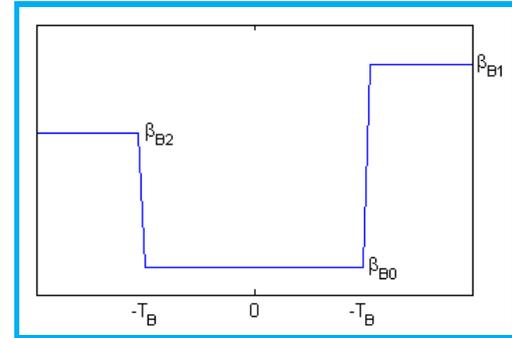


Magnified view of the temporal TIR and the evanescent wave associated with it.

- Fringe pattern is a consequence of interference between the incident and reflected pulses: Fringe spacing =  $\Delta t = (\Delta v)^{-1}$ .
- Temporal analog of the Goos-Hänchen shift also exists.  
Ponomarenko, Zhang, Agrawal, Phys. Rev. A **106** L061501 (2022).

# Temporal Waveguides

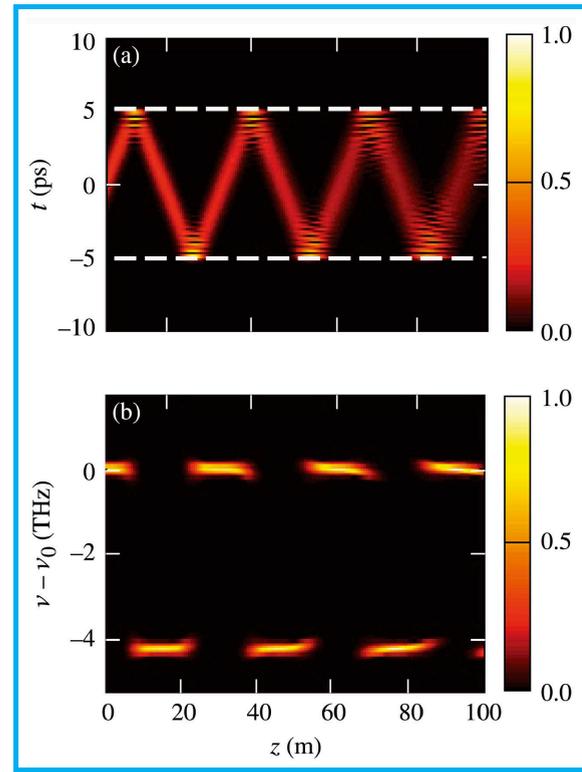
- TIR can be used to make temporal waveguides that trap optical pulses.
- Two temporal boundaries are needed.
- Refractive index of the central region can be lower or higher.



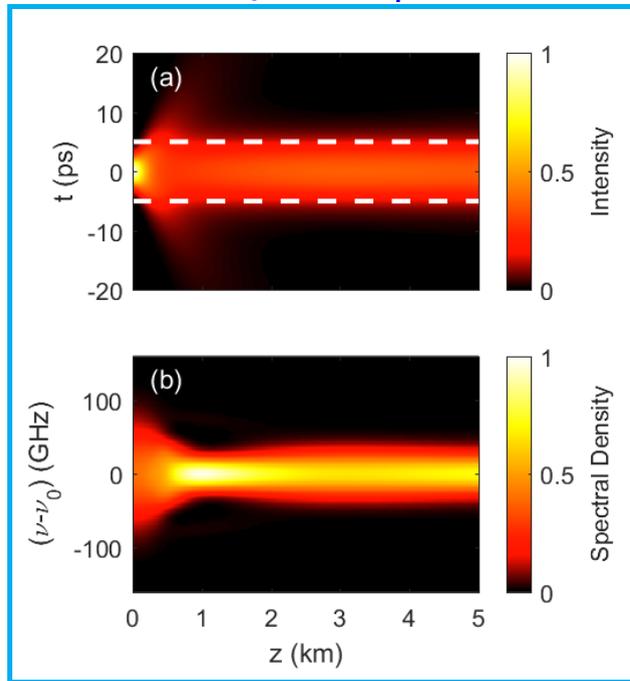
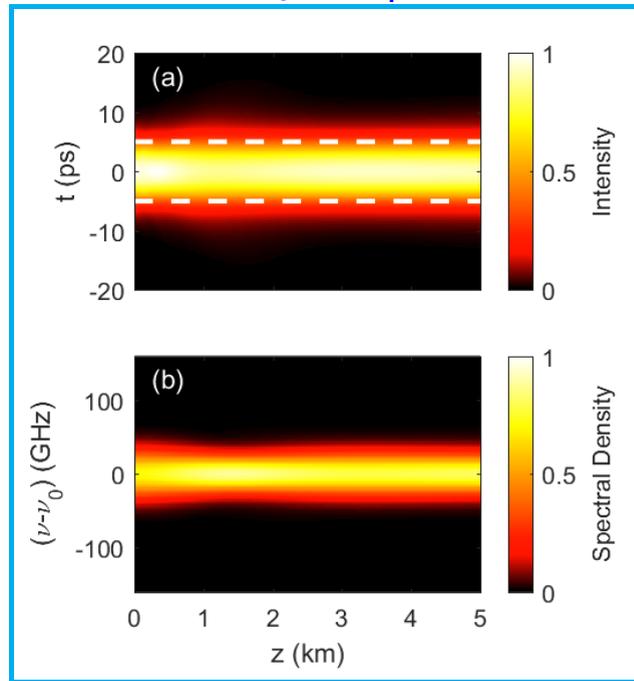
- A pulse is trapped inside the waveguide if it undergoes TIR at both temporal boundaries.
- Temporal waveguide support modes just as spatial waveguides do.
- The number of modes depends on the parameter  $V = 2\sqrt{k_0\Delta n T_B^2/\beta_2}$ .  
Plansinis, Donaldson, Agrawal, JOSA B **33**, 1112 (2016).

# Multimode Waveguides

- A 1-ps Gaussian pulse trapped inside a 10-ps wide waveguide:  $\beta_B = 5.6 \text{ m}^{-1}$ ,  $\Delta\beta_1 = 66.7 \text{ ps/km}$ ,  $\beta_2 = 50 \text{ ps}^2/\text{km}$
- Pulse undergoes TIR and its spectrum shifts after each reflection.
- Pulse broadening eventually creates distortions and pulse excites multiple waveguide modes ( $V = 26.8$ ).
- This approach can work when the propagation length is a fraction of the dispersion length ( $z < L_D$ ).



# Single-Mode Waveguides

 $T_0 = 2.5 \text{ ps}$  $T_0 = 5 \text{ ps}$ 

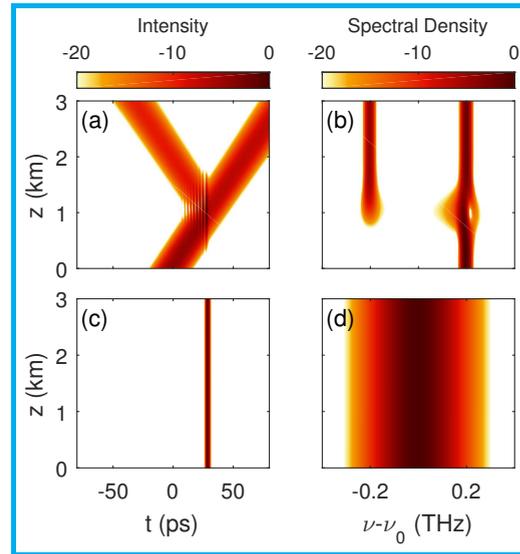
Plansinis, Donaldson, Agrawal, JOSA B **33**, 1112 (2016)

# Experimental Considerations

- Experimental realization of temporal waveguides requires two moving boundaries. Two techniques can be used for this purpose.
- A high-index temporal window can be created through the Pockels effect in an electro-optic medium using microwaves.
- An intense pump pulse can create a high-index region over its width in a nonlinear medium through the optical Kerr effect.
- Spreading of pump pulse can be controlled if its wavelength lies near the zero-dispersion wavelength of an optical fiber.
- This case has been studied numerically: Plansinis et al., J. Opt. Soc. Am. B **35**, 436 (2018).
- Simulations show that probe pulses can be confined to a temporal window created using super-Gaussian pump pulses.

# Temporal Reflection from a Soliton

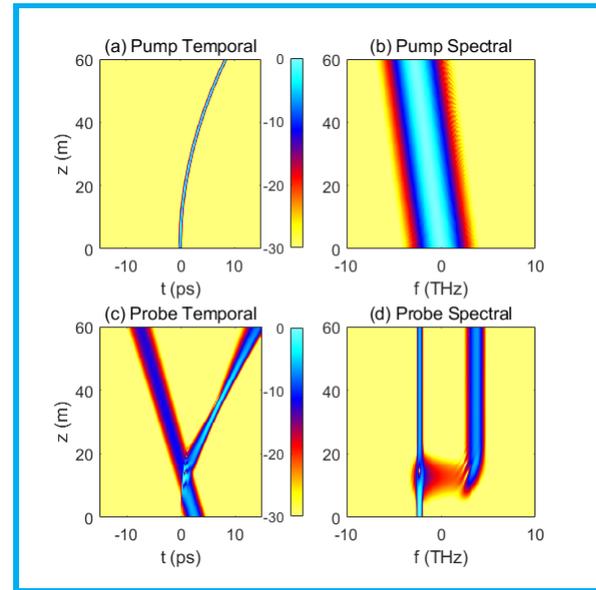
- Solitons are ideal for producing a moving temporal boundary.
- They require a medium with anomalous GVD ( $\beta_2 < 0$ ).
- Figure shows reflection of a probe pulse from a soliton: Plansinis et al., JOSA B **35**, 436 (2018).



- Pump pulse (3-ps wide) travels as a soliton inside the fiber ( $\beta_2 < 0$ ).
- Probe pulse (16-ps wide) arrives near soliton at 1 km ( $\beta_2 = 25 \text{ ps}^2/\text{km}$ ).
- It is partly reflected by the soliton with a large spectral shift.

# Impact of Raman scattering

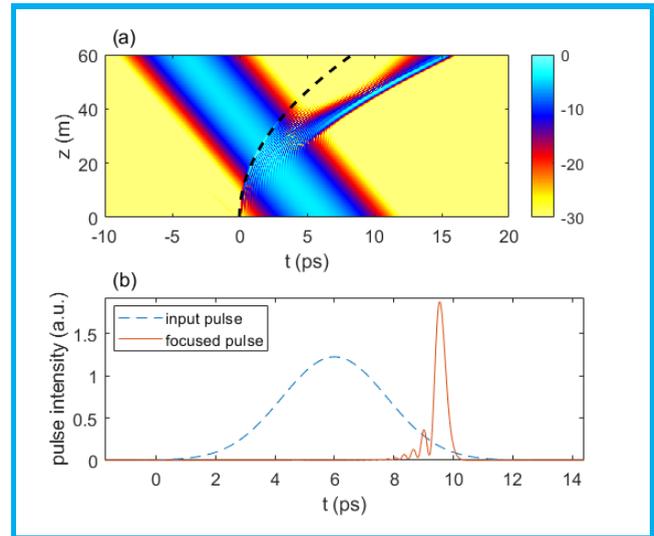
- Intrapulse Raman scattering (IRS) becomes relevant for short solitons.
- Solitons slow down as their spectrum red-shifts through IRS.
- Reflection of probe occurs from a curved temporal boundary.



- Reflection of a 4-ps probe pulse from a 100-fs soliton is shown.
- Soliton's trajectory is bent because of Raman-induced spectral changes.
- Zhang, Donaldson, Agrawal, JOSA B **39**, 1950 (2022).

# Temporal Focusing

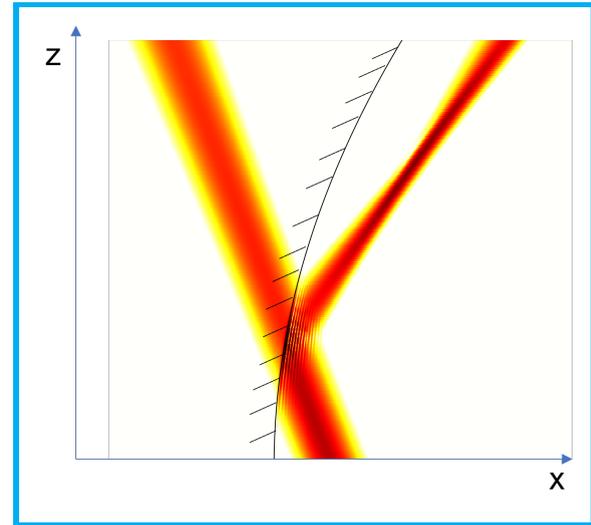
- Probe pulse exhibits temporal focusing after reflection.
- Its compression is induced by the curved soliton's trajectory acting as a temporal boundary.
- Zhang, Donaldson, and Agrawal, *JOSA B* **39**, 1950 (2022).



- Figure shows reflection of 4-ps probe pulse from 100-fs solitons.
- Reflected pulse is compressed by a factor of 10 at the focal point.
- The focusing is a consequence of the Raman-induced red-shift of the pump pulse traveling as a soliton.

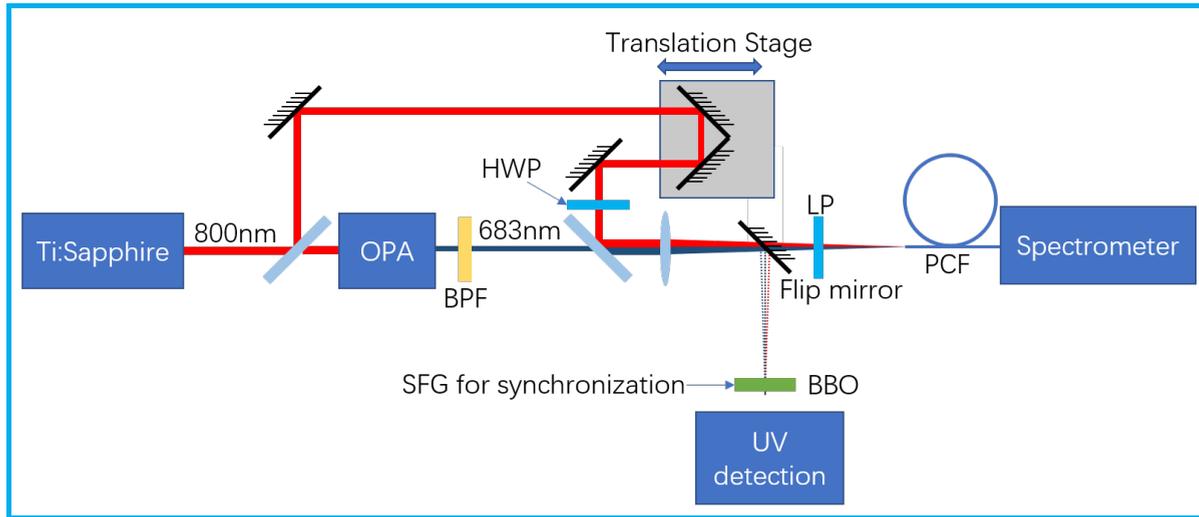
# Space-Time Duality for Focusing

- A parabolic mirror curves the phase front on reflection.
- This is the reason why the reflected light comes to a focus.
- Any optical Beam is compressed by a parabolic mirror at its focal point.



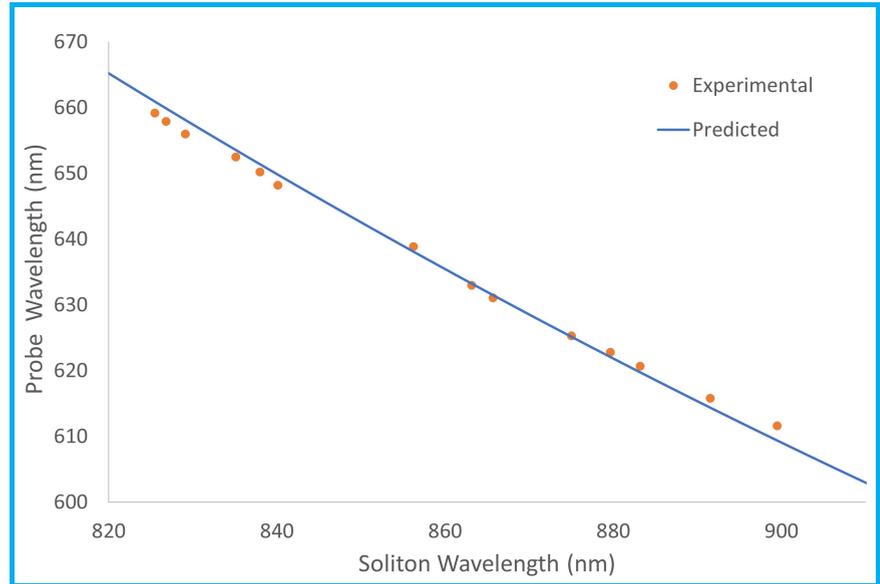
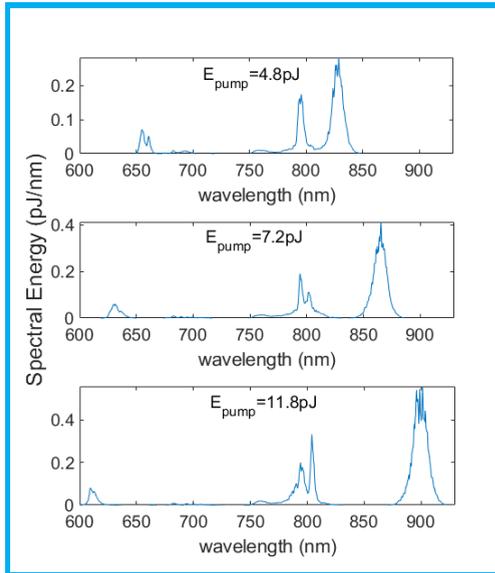
- Raman-induced temporal focusing inside a dispersive nonlinear medium is the temporal analog of a parabolic mirror.
- Curved trajectory of the soliton, resulting from frequency shifts induced by Raman scattering, mimics a parabolic mirror.

# Recent Experiment on Temporal Waveguides



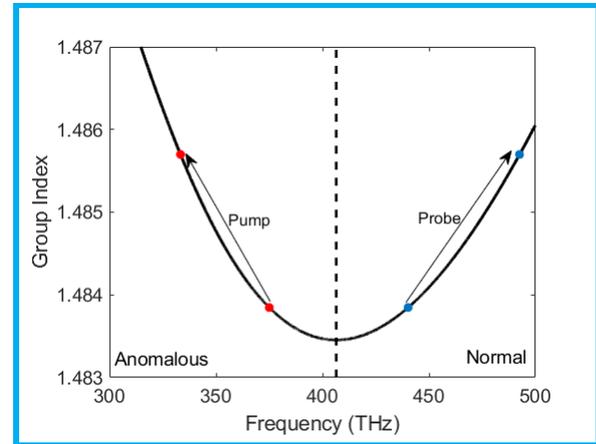
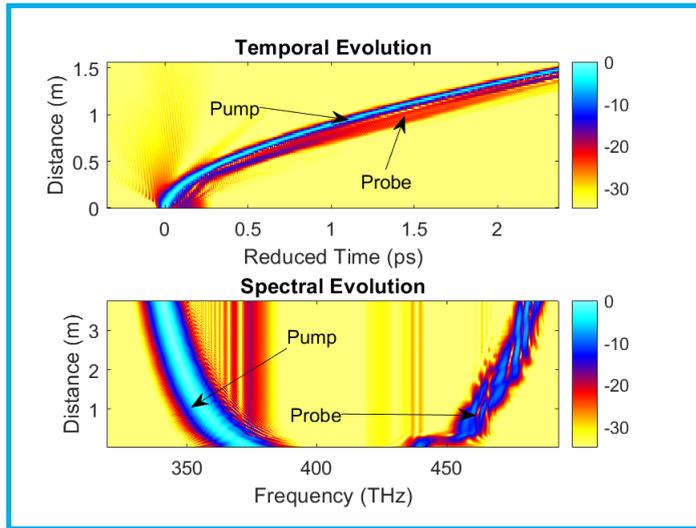
- A photonic crystal fiber (3.75-m-long) used as a dispersive nonlinear medium.
- Pump pulses were only 30-fs wide at the 800-nm wavelength.
- Probe pulses at 683-nm were  $> 1$  ps wide after the bandpass filter.
- Output spectra recorded for different energies of pump pulses,

# Measured Output Spectra



- Pump wavelength shifts toward red because of Raman scattering.
- Probe wavelength shifts toward blue because of temporal waveguiding.
- These shifts agree with predictions based on dispersion of the fiber.

# Simulations and Dispersion Data



- Probe follows pump's trajectory because of multiple temporal reflections.
- The two pulses shift their spectra on opposite sides to ensure that they move at the same speed.
- Measured spectral shifts agree fully with the dispersion data.

# Conclusions

- Inclusion of chromatic dispersion is relevant for pulses propagating inside a time-varying medium.
- Kerr nonlinearity of a dispersive medium can be exploited to create moving temporal boundaries using intense pump pulses.
- Temporal equivalent of total internal reflection occurs for optical pulses at a moving temporal boundary.
- A waveguide is formed by two such temporal boundaries.
- Such waveguides are temporal analog of planar optical waveguides.
- Raman scattering leads to new effects when femtosecond pulses are propagated as solitons.
- An experiment has shown that such short solitons, forming inside a photonic crystal fiber, can guide probe pulses along their trajectories.