



Nonlinear Performance of SDM Systems Designed with Multimode or Multicore Fibers

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Introduction

- Internet traffic has been growing exponentially every year.
- System capacity must follow this growth to avoid a capacity crunch.
- The use of WDM, PSK, and PDM helped us during the last decade.
- Space-division multiplexing (SDM) may help us during this decade.



(Essiambre and Tkach, IEEE Proc., May 2012)





Space-division multiplexing







System Design Issues

- Linear coupling among fiber modes would invariably lead to channel crosstalk.
- Nonlinear coupling among fiber modes will also occur through cross-phase modulation (XPM) an four-wave mixing (FWM).
- How does the linear coupling affect nonlinear penalties in a specific channel of a SDM system?
- If polarization-division-multiplexing (PDM) is employed in addition to WDM and SDM, what is the impact of birefringence fluctuations?
- To answer such questions, it is necessary to develop a comprehensive numerical model for SDM systems.
- In this talk, I present our theoretical model and recent numerical results. Mumtaz et al., PTL 24, 1574-1576 (2012); JLT 31, 398-406 (2013).



Theoretical Framework

• We start with the wave equation in the frequency domain

$$\nabla^{2}\widetilde{\mathbf{E}}(\boldsymbol{\omega}) + n^{2}(x,y)\frac{\boldsymbol{\omega}^{2}}{c^{2}}\widetilde{\mathbf{E}}(\boldsymbol{\omega}) = -\boldsymbol{\omega}^{2}\boldsymbol{\mu}_{0}\widetilde{\mathbf{P}}^{\mathrm{NL}}(\boldsymbol{\omega}).$$

• Nonlinear effects included in the time domain through

$$\mathbf{P}^{\mathrm{NL}}(t) = \frac{\boldsymbol{\varepsilon}_0}{4} \boldsymbol{\chi}^{(3)} \left[(\mathbf{E} \cdot \mathbf{E}) \mathbf{E}^* + 2(\mathbf{E}^* \cdot \mathbf{E}) \mathbf{E} \right].$$

• Expand the optical field in terms of fiber modes

$$\widetilde{\mathbf{E}}(\mathbf{r},\boldsymbol{\omega}) = \sum_{m=1}^{M} \mathbf{F}_m(x,y) \widetilde{\mathbf{A}}_m(z,\boldsymbol{\omega}) e^{i\boldsymbol{\beta}_m(\boldsymbol{\omega})z},$$

• $\widetilde{\mathbf{A}}_m = [\widetilde{\mathbf{A}}_{mx} \ \widetilde{\mathbf{A}}_{my}]^{\mathrm{T}}$ includes both polarization components of the *m*th fiber mode (needed for PDM).

Nonlinear Propagation Equations

• The time-domain equation for the *p*th mode satisfies

$$\frac{\partial \mathbf{A}_p}{\partial z} - i(\boldsymbol{\beta}_{0p} - \boldsymbol{\beta}_r)\mathbf{A}_p + \left(\boldsymbol{\beta}_{1p} - \frac{1}{v_{g_r}}\right)\frac{\partial \mathbf{A}_p}{\partial t} + \frac{i\boldsymbol{\beta}_{2p}}{2}\frac{\partial^2 \mathbf{A}_p}{\partial t^2}$$
$$= i\sum_m q_{mp}\mathbf{A}_m + \frac{i\gamma}{3}\sum_{lmn} f_{lmnp}\left[(\mathbf{A}_n^{\mathrm{T}}\mathbf{A}_m)\mathbf{A}_l^* + 2(\mathbf{A}_l^{\mathrm{H}}\mathbf{A}_m)\mathbf{A}_n\right].$$

• Linear and nonlinear couplings among spatial modes governed by

$$q_{mp} \propto k_0 \iint \Delta n(x, y, z) \mathbf{F}_p^* \mathbf{F}_m dx dy, \quad f_{lmnp} = \iint \mathbf{F}_l^* \mathbf{F}_m \mathbf{F}_n \mathbf{F}_p^* dx dy.$$

• Fiber modes are normalized such that

$$\iint \mathbf{F}_p^*(x,y)\mathbf{F}_m(x,y)dx\,dy = \boldsymbol{\delta}_{mp}$$

Linear Coupling in Multimode Fibers

- Coupling between modes does not occur in a perfect multimode fiber.
- In practice, spatial modes experience random coupling owing to fiber imperfections.
- Coupling is governed by small fluctuations in the refractive index $[\Delta n(x, y, z) \sim 10^{-4}].$

This case will be discussed later in this talk.

Linear Coupling in Multicore Fibers

- Coupling between two cores depends on the spacing between them.
- Coupling coefficient q_{mp} is a measure of this coupling.
- Coupling length $L_c = \pi/(2q)$ can vary from <1 cm to >1000 km depending on how close the two cores are.

Nonlinear Coupling in Multicore Fibers

- Governed by the mode-overlap factor $f_{lmnp} \propto \iint F_l^* F_m F_n F_p^* dx dy$.
- Self-phase modulation (SPM) occurs through $f_{pppp} = 1$.
- Cross-phase modulation (XPM) governed by terms like f_{ppmm} .
- Other Combinations represent FWM-like effects.

XPM and FWM terms are negligible in multicore fibers whose cores of radius a are separated by d > 4a.

Numerical Results for Multicore Fibers

- We considered multicore fibers with up to 19 cores.
- Two 28.5-Gbaud PDM-QPSK symbol streams launched into each core.
- Propagation over 1000 km by the split-step Fourier method.
- Least square equalizer (LSE) is based on a training sequence.
- Actual bit error rate (BER) calculated using 2^{20} bits during simulations.

Three-core Fibers

Stronger linear coupling (shorter L_c) improves the BER at a given OSNR.

OSNR penalties

OSNR penalties calculated after 1000 km using $L_c = 100$ m.

SPM-Induced Noise

- Nonlinear noise after 100 km for an input power of 12 dBm.
- Degradations induced by SPM are reduced in a 3-core fiber.

Multicore Fibers

Simulated BER curves for fibers containing up to 7 cores.

Multicore Fibers (cont.)

Simulated BER curves for fibers containing up to 19 cores.

Birefringence Fluctuations

- Birefringence fluctuations must be included in realistic simulations.
- Their inclusion makes all mode-propagation equations stochastic.
- Numerical simulations become time-consuming as one must solve these equations many times before averaging.
- In the case of single-mode fibers, Manakov equations are found to be quite useful for predicting the average behavior.
- We have derived new Manakov equations for multimode fibers assuming strong polarization coupling but no coupling among spatial modes.
- Fiber divided into many sections with randomly oriented phase plates [Mumtaz et al., JLT **31**, 398 (2013)].

Manakov Equations for Multimode Fibers

• After averaging over birefringence fluctuations, we obtain

$$\frac{\partial \bar{\mathbf{A}}_p}{\partial z} + d_p \frac{\partial \bar{\mathbf{A}}_p}{\partial t} + \frac{i\beta_{2p}}{2} \frac{\partial^2 \bar{\mathbf{A}}_p}{\partial t^2} = i\gamma \Big(f_{pppp} \frac{8}{9} |\bar{\mathbf{A}}_p|^2 + \sum_{m \neq p} f_{mmpp} \frac{4}{3} |\bar{\mathbf{A}}_m|^2 \Big) \bar{\mathbf{A}}_p.$$

- Modal group-velocity mismatch included through d_p .
- The SPM term contains the usual factor of 8/9 as in single-mode fibers.
- The factor of 2 in the XPM terms is reduced to 4/3.
- This reduction indicates that XPM penalties are reduced in multimode fibers in the presence of birefringence fluctuations.
- The contribution of all FWM terms averages out to zero.
- Numerical simulation support these conclusions.

Numerical Results for Multimode Fibers

- We consider both step-index and graded-index multimode fibers.
- Core diameter 12 μ m, $\Delta = 0.01$, V = 5 at 1550 nm.

Spatial distribution shapes	Step-index fiber	Graded-index fiber
	LP01	HG00
	LP11a, LP11b	HG01, HG10
	LP02	HG02+HG20
	LP21a, LP21b	HG11a, HG11b

Numerical Results for Multimode Fibers

- PDM-QPSK symbol streams launched into each mode (bit rate 114 G/s).
- Propagation over 1000 km by the split-step Fourier method.
- Solid curves, Manakov; Squares, full model; circles, no birefringence.

Manakov Equations in Strong-Coupling Regime

- Coupling between spatial modes can occur in practice.
- In general, coupling strength varies for various mode pairs and is strongest for nearly degenerate modes.
- Strong-coupling regime (all modes strongly coupled) is easy to deal with [(Mecozzi et al., Opt. Exp. **20**, 11673 (2012)].
- Manakov equation in this regime takes the form

$$\frac{\partial \bar{\mathscr{A}}}{\partial z} + \frac{i\bar{\beta}_2}{2}\frac{\partial^2 \bar{\mathscr{A}}}{\partial t^2} = i\gamma\kappa|\bar{\mathscr{A}}|^2\bar{\mathscr{A}}, \quad \kappa = \sum_{k=1}^M \sum_{l\leq k}^M \frac{32}{2^{\delta_{kl}}}\frac{f_{kkll}}{6M(2M+1)}.$$

• Here $\overline{\mathscr{A}}$ is a column vector with 2M elements (M modes, 2 polarizations).

Numerical Results

- PDM-QPSK symbol streams launched into each mode (bit rate 114 Gb/s).
- Propagation over 1000 km by the split-step Fourier method.
- Coupling regimes compared in the case of a graded index fiber.

Mode Coupling in Multimode Fibers

- Neither of the two preceding coupling models is realistic.
- Spatial modes of an ideal fiber are not coupled. In this case, two orthogonally polarized components of each mode are coupled through birefringence fluctuations (no coupling among spatial modes).
- In practice, random fluctuations in the refractive index of a multimode fibre result in coupling of all spatial modes.
- However, not all modes are strongly coupled since coupling also depends on how close the propagation constants of the two modes are.
- A more general coupling model is needed that takes into account such mode-dependent coupling.

General Mode-Coupling Model

- We are developing a model that takes into account of the fact that the strength of mode coupling varies from one mode pair to another.
- To account for random refractive index fluctuations, we assume $\mathcal{E}(x, y, z) = n^2(x, y) + \Delta \mathcal{E}(x, y, z)$, where $\Delta \mathcal{E}(x, y, z)$ is a random variable.
- $\Delta \varepsilon(x, y, z)$ vanishes on average at every point in the transverse plane such that

$$\langle \Delta \varepsilon(x, y, z) \Delta \varepsilon(x, y, z') \rangle = \sigma_{\varepsilon}^2 \exp(-|z - z'|/l_d)$$

Coupling and Transfer Matrices

- Fluctuations $\Delta \varepsilon(x, y, z)$ couple all spatial modes of a multimode fiber through a random coupling matrix.
- Average value of all elements of this martix is zero but the variance depends on the standard deviation of index fluctuations.
- A transfer matrix can be used to propagate all modes simultaneously along the fiber.
- It turn out that the strength of coupling between two modes is governed by the ratio $q_{mn} = \kappa_{mn}/|\beta_m \beta_n|$.
- Here κ_{mn} is the root-mean square value of the coupling coefficient and β_m is the propagation constant of the *m*th mode.
- As expected, modes with smaller values of $|\beta_m \beta_n|$ exhibit stronger coupling.

Concluding Remarks

- The use of SDM through multimode or multicore fibers is promising for future telecom systems.
- We have studied the impact of mode coupling on nonlinear penalties in such systems.
- In a multicore fiber, nonlinear penalties are reduced when coupling is enhanced by bringing cores together.
- Random birefringence must be included in all cases.
- Two new Manakov equations derived in the weak and strong coupling regimes by averaging over random birefringence.
- A new coupling model is under development. It will allow us to study the transition from no coupling to strong coupling.